

# Why quantum logic cannot be classical

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# Classical event structure

**$\sigma$ -algebra** of sets,  $\mathcal{L} \subseteq 2^U$ :

- $U \in \mathcal{L}$
- $\{A_i \mid i \in \mathbb{N}\} \subseteq \mathcal{L} \implies \bigcup_{i \in \mathbb{N}} A_i \in \mathcal{L}$

**State** (=probability measure)  $s: \mathcal{L} \rightarrow [0, 1]$ :

- $s(U) = 1$
- $\{A_i \mid i \in \mathbb{N}\} \subseteq \mathcal{L}, A_i \cap A_j = \emptyset \text{ for } i \neq j \implies s\left(\bigcup_{i \in \mathbb{N}} A_i\right) = \sum_{i \in \mathbb{N}} s(A_i)$

$\mathcal{S}(\mathcal{L}) :=$  **state space** of  $\mathcal{L}$ ; it is a Choquet simplex

**Pure states:** extreme points of  $\mathcal{S}(\mathcal{L})$

**Two-valued states:**  $\mathcal{S}(\mathcal{L}) \cap \{0, 1\}^{\mathcal{L}}$

For  $\sigma$ -algebras:

- pure states = two-valued states = points in the Stone space
- state space (even the space of two-valued states) determines the whole structure

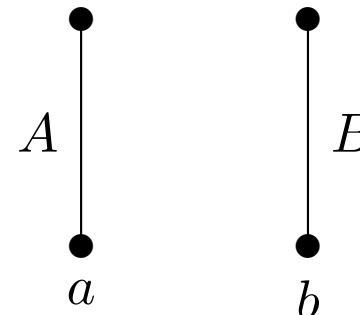
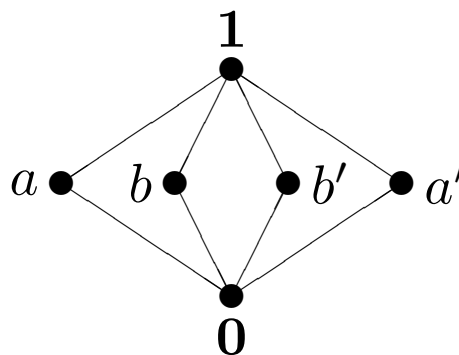
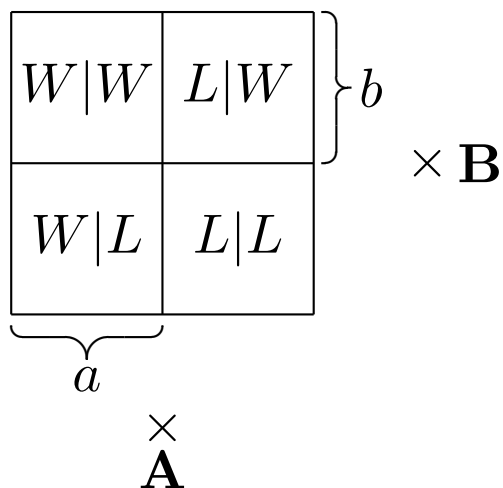
We need **disjoint**, not all unions!

**$\sigma$ -class** of sets,  $\mathcal{L} \subseteq 2^U$ :

- $U \in \mathcal{L}$
- $\{A_i \mid i \in \mathbb{N}\} \subseteq \mathcal{L}, A_i \cap A_j = \emptyset \text{ for } i \neq j \implies \bigcup_{i \in \mathbb{N}} A_i \in \mathcal{L}$

# Example 1

$W/L = \text{Wins/Loses with|without player JJ}$



$$U = \{W|W, W|L, L|W, L|L\}$$

$$A = \{\emptyset, \underbrace{\{W|W, W|L\}}_a, \underbrace{\{L|W, L|L\}}_{a'}, U\}$$

$$B = \{\emptyset, \underbrace{\{W|W, L|W\}}_b, \underbrace{\{W|L, L|L\}}_{b'}, U\}$$

$$\mathcal{L} = A \cup B = \{\emptyset, a, a', b, b', U\}$$

$\mathcal{L}$  is a (nondistributive modular) lattice called *MO2*

Pure states:

$$s : \mathcal{L} \rightarrow \{0, 1\}$$

$s(A)$	$s(B)$
0	0
0	1
1	0
1	1

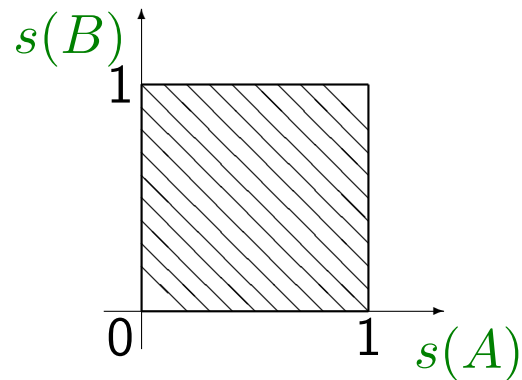
(S0)  $s(U) = 1$

(S1)  $s(x') = 1 - s(x)$

All states:

$s : \mathcal{L} \rightarrow [0, 1]$ , satisfy (S0), (S1)

$s(A) = p, s(B) = q, p, q \in [0, 1]$  arbitrary



## Example 2

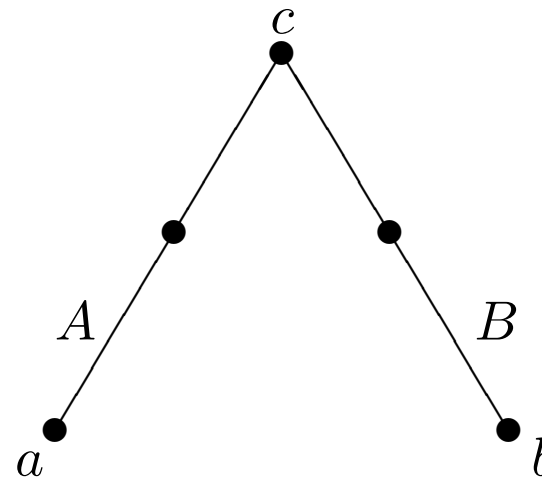
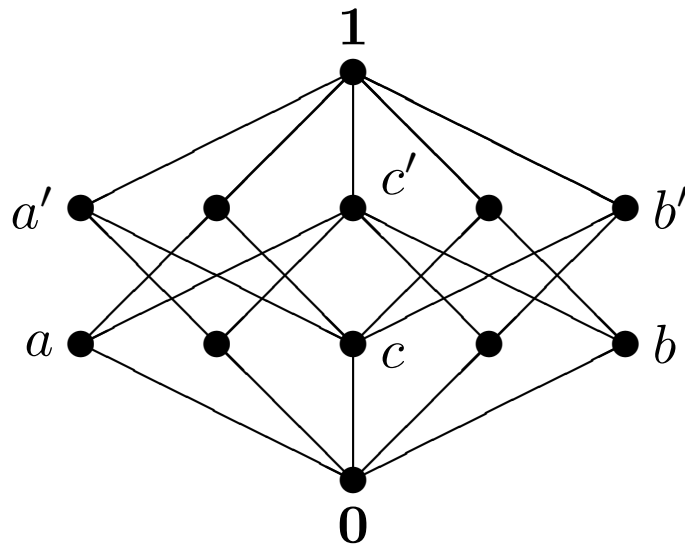
Example 1 with one more result,  $c = \text{match cancelled}$

$$A = \{0, a, c, (a \vee c)', a \vee c, a', c', 1\}$$

$$B = \{0, b, c, (b \vee c)', b \vee c, b', c', 1\}$$

$$A \cap B = \{0, c, c', 1\}$$

$$\mathcal{L} = A \cup B = \{0, a, b, c, a \vee c, b \vee c, (a \vee c)', (b \vee c)', a', b', c', 1\}$$

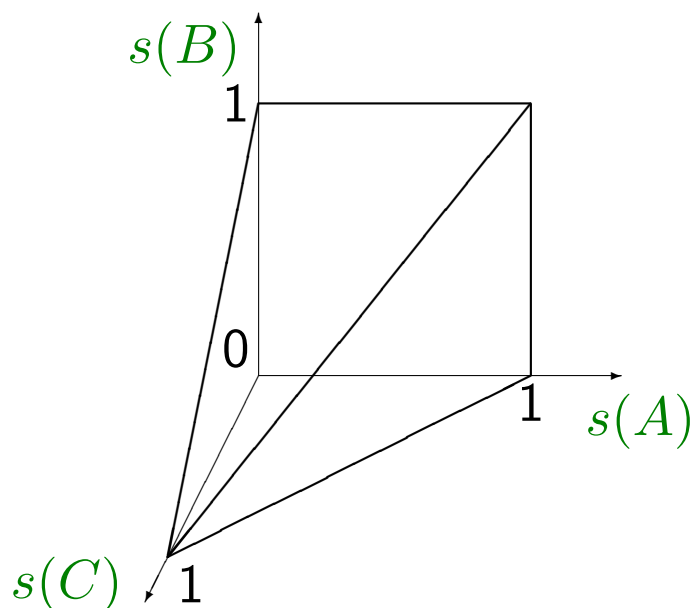


Pure states:

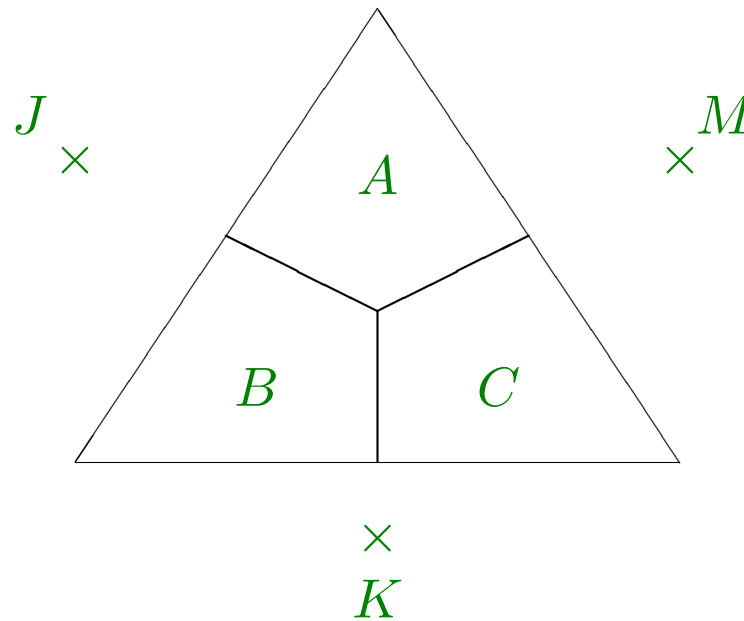
$s(A)$	$s(B)$	$s(C)$
0	0	0
0	1	0
1	0	0
1	1	0
0	0	1

All states:

$$s(A) = p, \quad s(B) = q, \quad s(C) = r, \quad r \in [0, 1] \text{ arbitrary, } p, q \in [0, 1 - r]$$

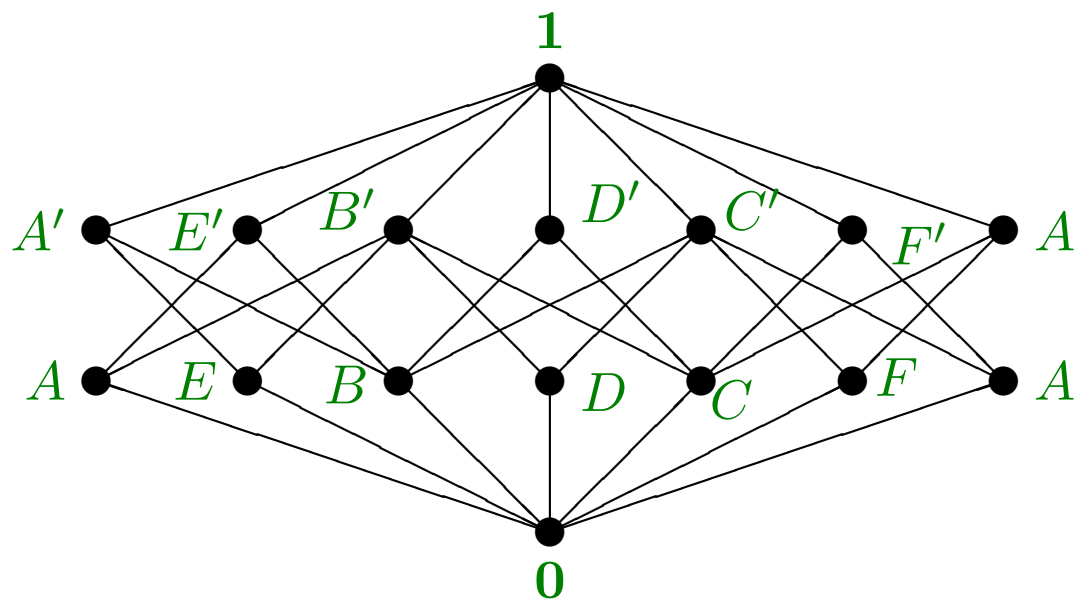
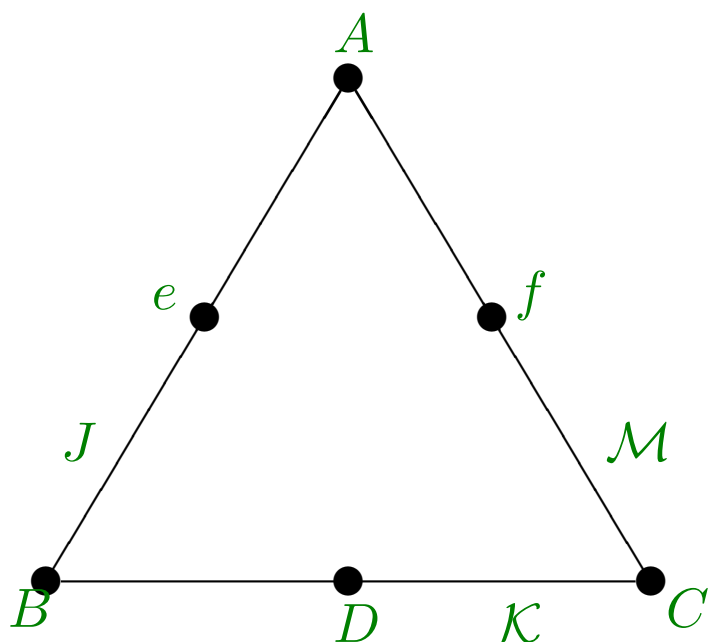


### Example 3:



$\mathcal{K} = \{0, B, C, D, B', C', D', 1\}$ , where  $D$  means “the fire-fly is not observed from  $K$ ”  
 $\mathcal{M} = \{0, A, C, E, A', C', E', 1\}$ , where  $A$  means “the fire-fly is observed in the upper part”  
 $\mathcal{J} = \{0, A, B, F, A', B', F', 1\}$   
 $\mathcal{K} \cup \mathcal{M} \cup \mathcal{J} = \{0, A, B, C, D, E, F, A', B', C', D', E', F', 1\}$

This is **not** a lattice.



Pure states:

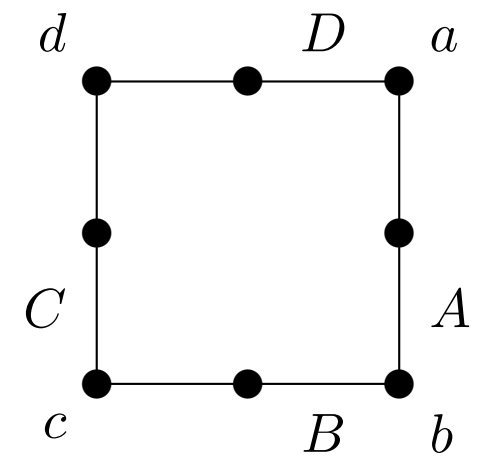
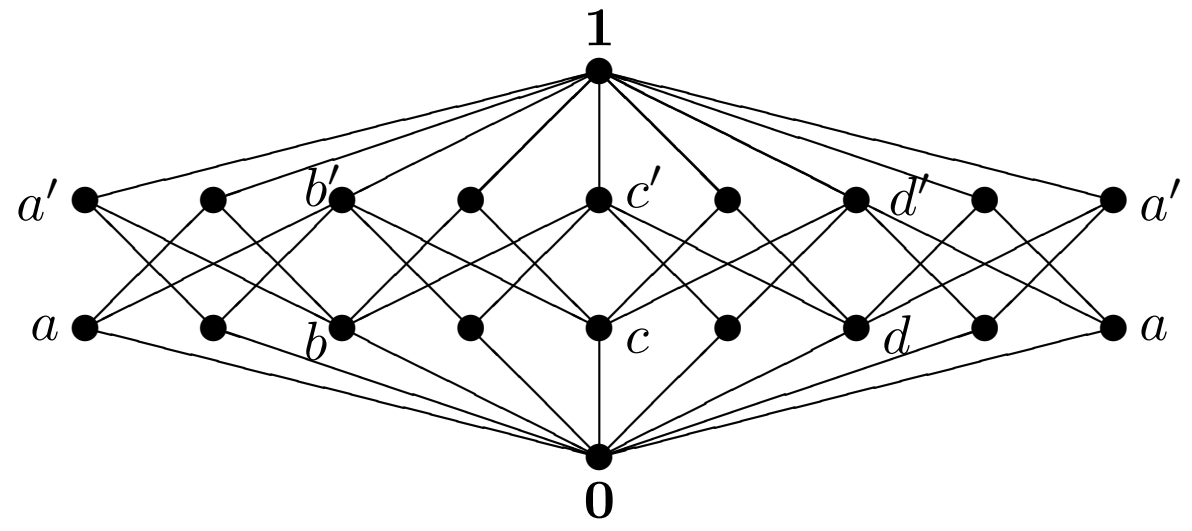
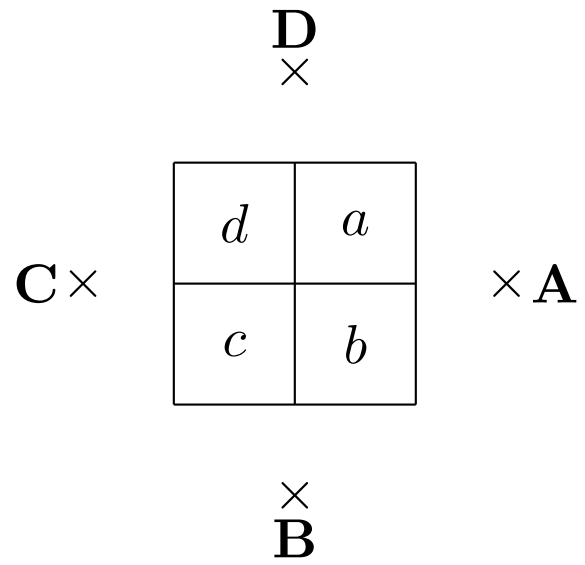
$s(A)$	$s(B)$	$s(C)$
1	0	0
0	1	0
0	0	1
0	0	0
1/2	1/2	1/2

All states:

$$s(A) = p, s(B) = q, s(C) = r, \quad p, q, r \in [0, 1], \quad p + q \leq 1, \quad p + r \leq 1, \quad q + r \leq 1$$



# Example 3 (non-transparent barriers)



# Example 3 (non-transparent barriers)



Pure states:

$s(a)$	$s(b)$	$s(c)$	$s(d)$
1	0	1	0
0	1	0	1
1	0	0	0
0	1	0	0
0	0	1	0
0	0	0	1
0	0	0	0

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## Orthomodular lattices

*Definition:* An **orthomodular lattice** is a lattice with bounds  $\mathbf{0}$ ,  $\mathbf{1}$  equipped with a unary operation  $' : \mathcal{L} \rightarrow \mathcal{L}$  (**orthocomplementation**) such that, for all  $a, b \in \mathcal{L}$ ,

- $a'' = a$
- $a \leq b \implies b' \leq a'$
- $a \wedge a' = \mathbf{0}$
- $a \leq b \implies b = a \vee (a' \wedge b)$  (**orthomodular law**)

**Orthogonality:**  $a \perp b \iff a \leq b'$

(This condition is strictly stronger than the usual  $a \wedge b = \mathbf{0}$ .)

*Example:* A  $\sigma$ -class of subsets needs not be a lattice; if it is, it is an OML.

## Structure of orthomodular lattices

**Boolean subalgebra:**  $\mathcal{M} \subseteq \mathcal{L}$  such that

- $0, 1 \in \mathcal{M}$ ,
- $a \in \mathcal{M} \implies a' \in \mathcal{M}$ ,
- $(\mathcal{M}, \leq|_{\mathcal{M}}, ' |_{\mathcal{M}})$  is a Boolean algebra.

**Compatibility:**  $a \leftrightarrow b \iff \exists$  Boolean subalgebra  $\mathcal{M}: a, b \in \mathcal{M}$

**Block:** a maximal Boolean subalgebra

**Center:** The set of all  $a \in \mathcal{L}$  such that  $\forall b \in \mathcal{L} : a \leftrightarrow b$

= the set of all “absolutely compatible” elements

= the classical part of the system

= the intersection of all blocks

**Atom:**  $a \in \mathcal{L} \setminus \{0\}$  such that there is no  $b$  satisfying  $0 < b < a$

$\mathcal{A}(\mathcal{L}) :=$  the set of all atoms of  $\mathcal{L}$

( $\sigma$ -additive) **state:**  $s: \mathcal{L} \rightarrow [0, 1]$  such that

- $s(1) = 1$
- $\{a_i \mid i \in \mathbb{N}\} \subseteq \mathcal{L}, a_i \perp a_j$  for  $i \neq j \implies s\left(\bigvee_{i \in \mathbb{N}} a_i\right) = \sum_{i \in \mathbb{N}} s(a_i)$

# Orthomodular lattices as families of Boolean algebras

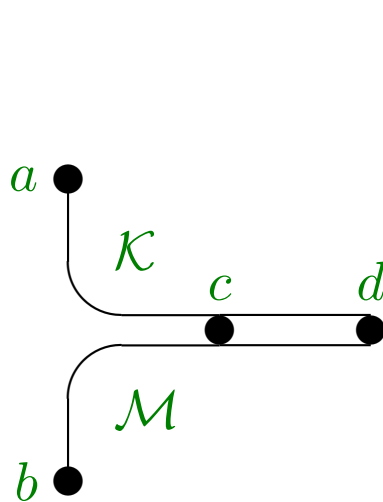
Every OML is the union of its maximal Boolean subalgebras (=blocks)

**Hypergraph:** a nonempty set (of **vertices**) and its covering by nonempty subsets (**edges**)

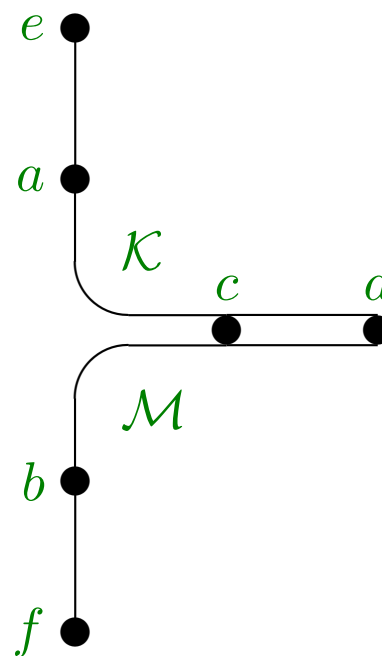
**Greechie diagram:** hypergraph whose vertices are atoms and edges are blocks

**State on a hypergraph:** evaluation of vertices such that the sum over each edge is 1

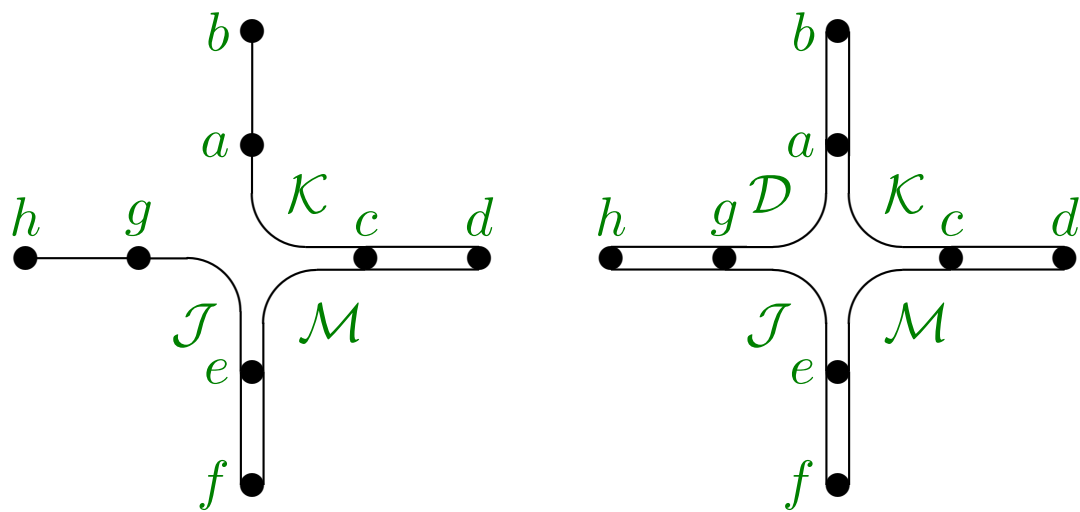
*Problem:* Which hypergraphs are Greechie diagrams of OMLs?



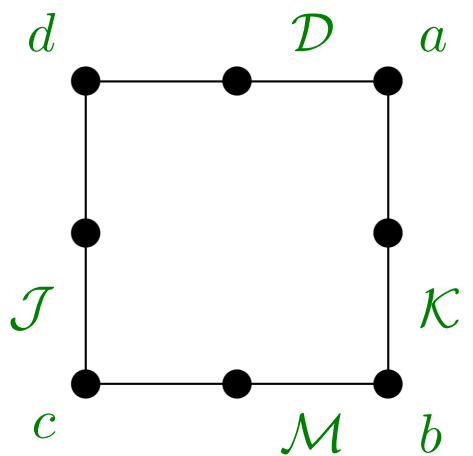
$$a = (c \vee d)' = b$$



$$a \vee e = b \vee f$$

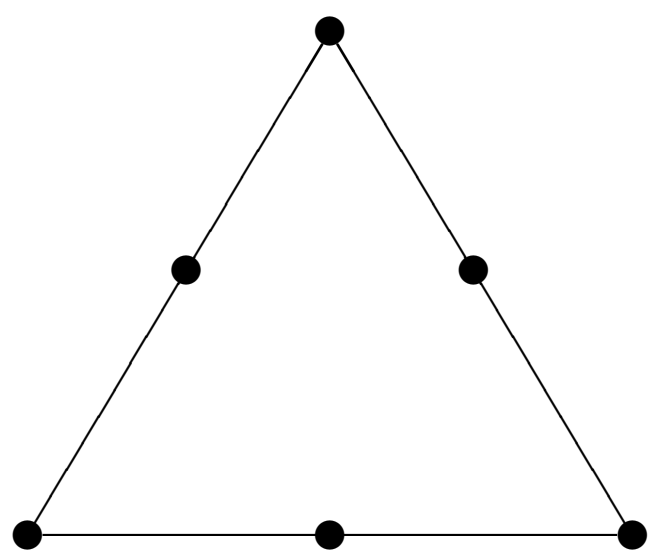


$$a \vee b = e \vee f \perp g \vee h$$

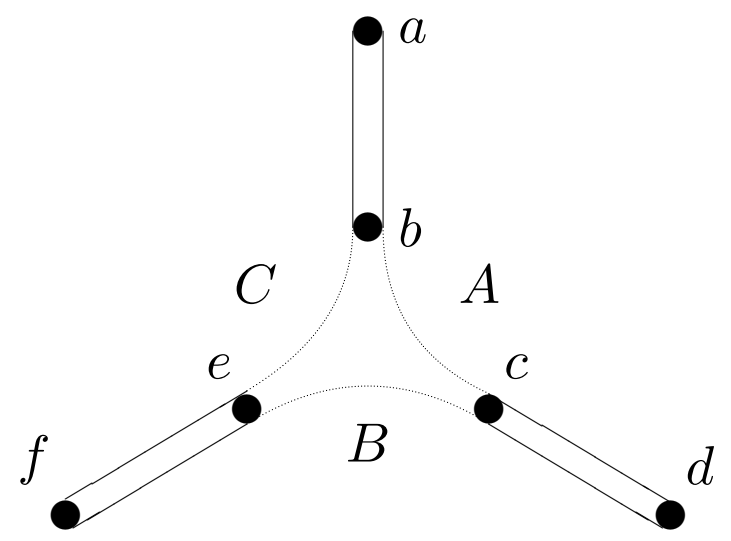


# Orthoalgebras

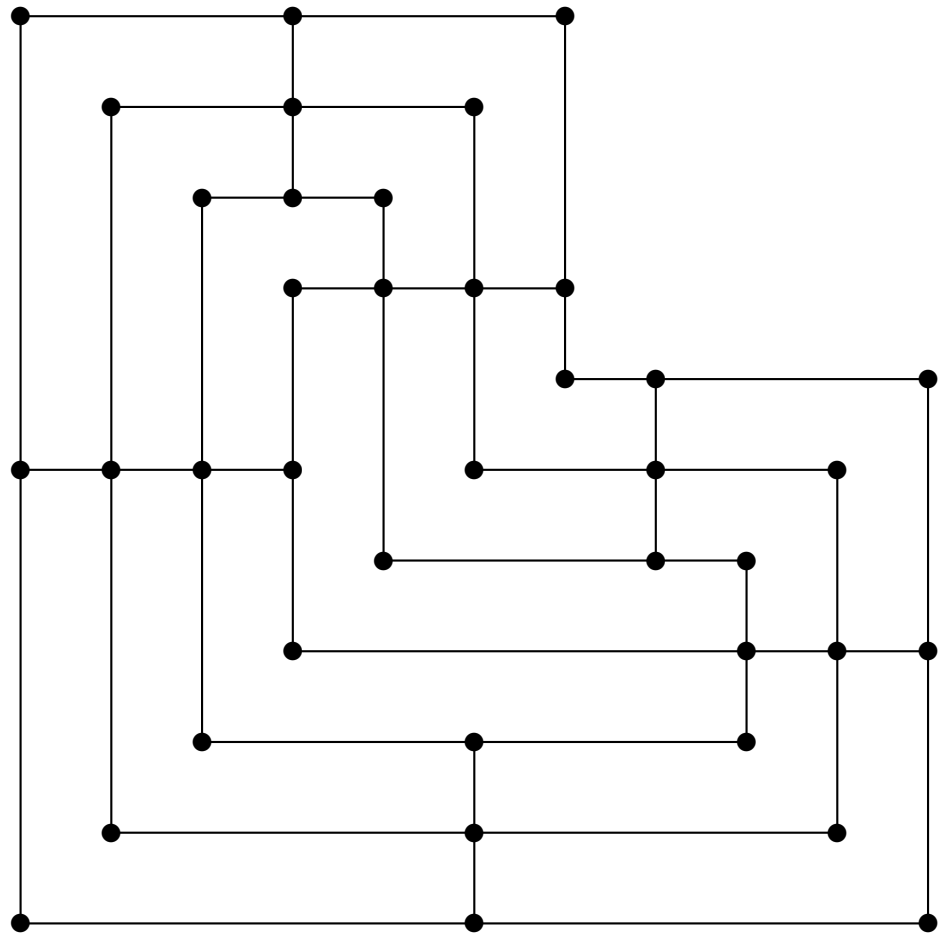
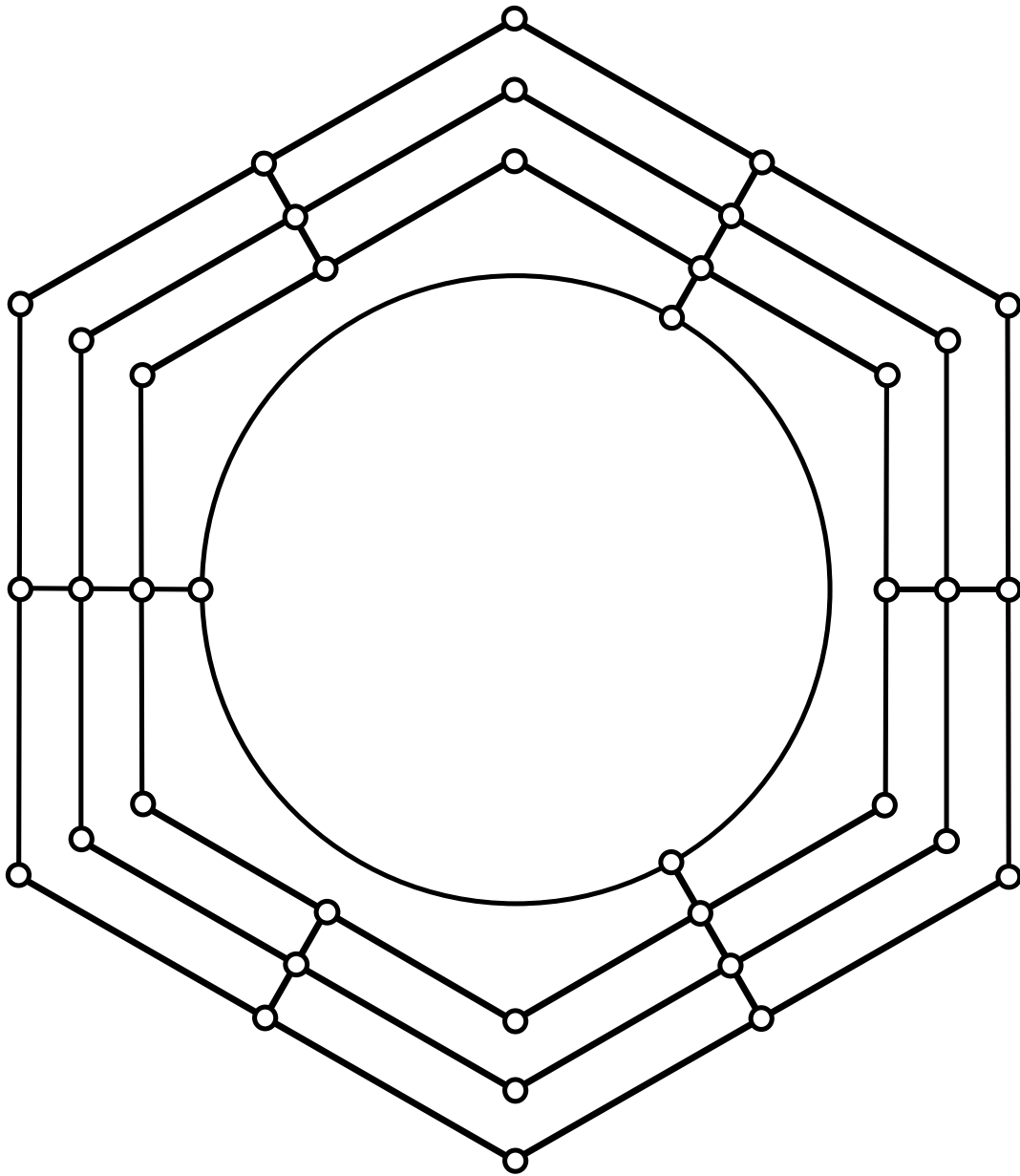
Allowed



Forbidden

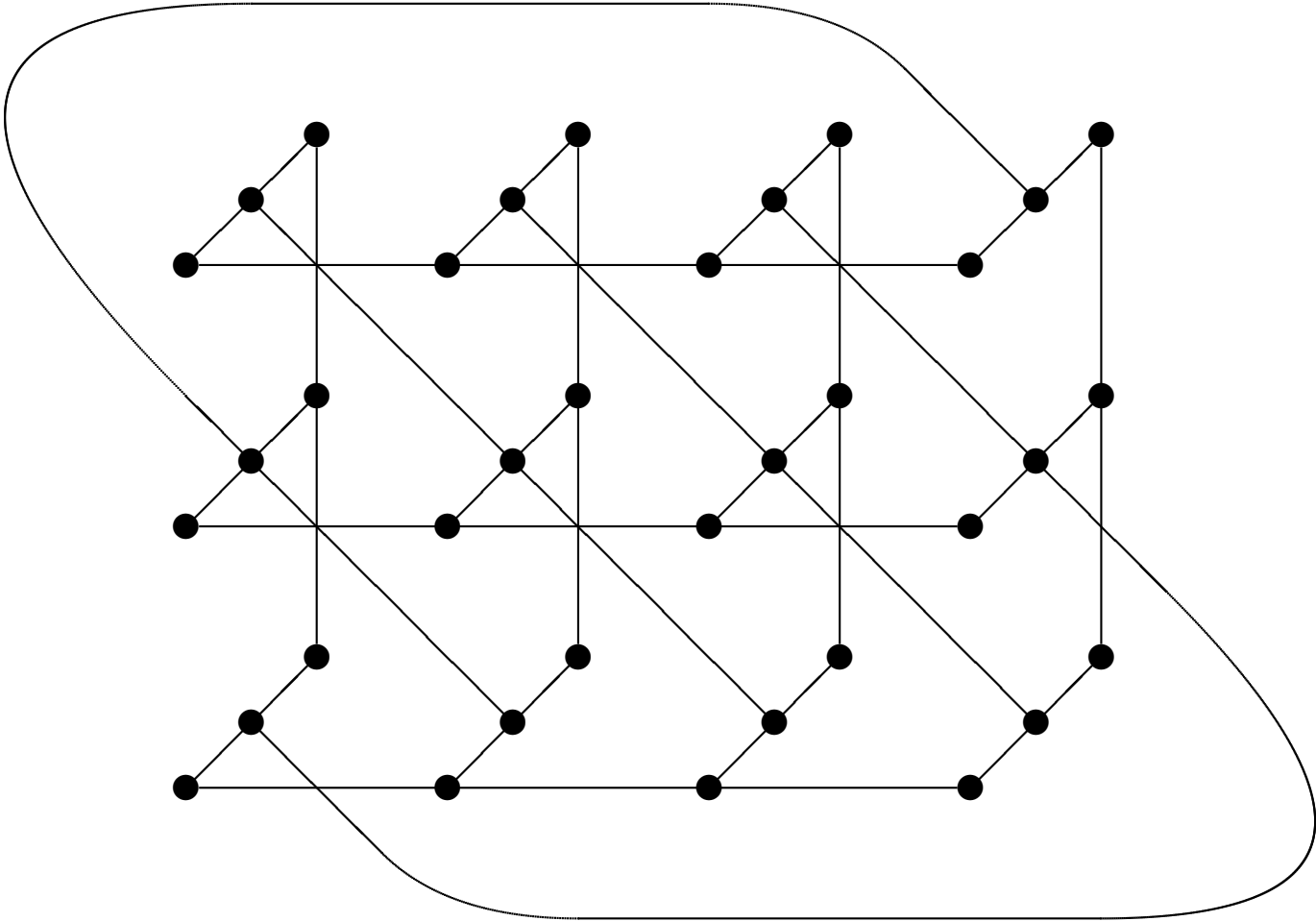


In particular, the state space may be empty [Rogalewicz]:

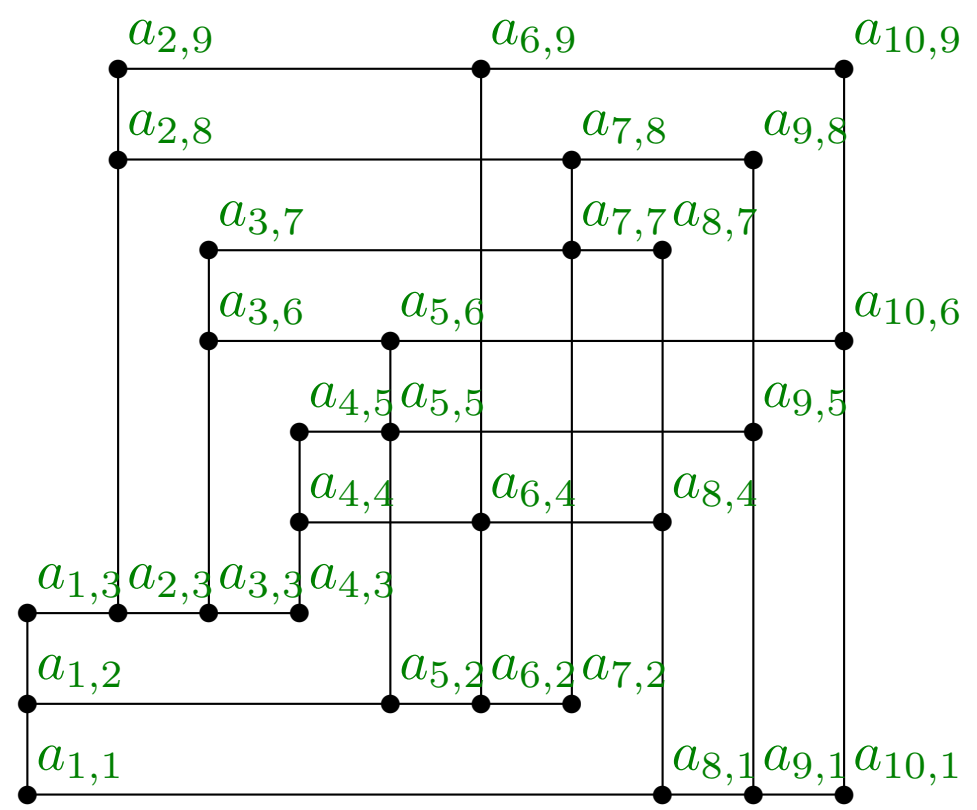
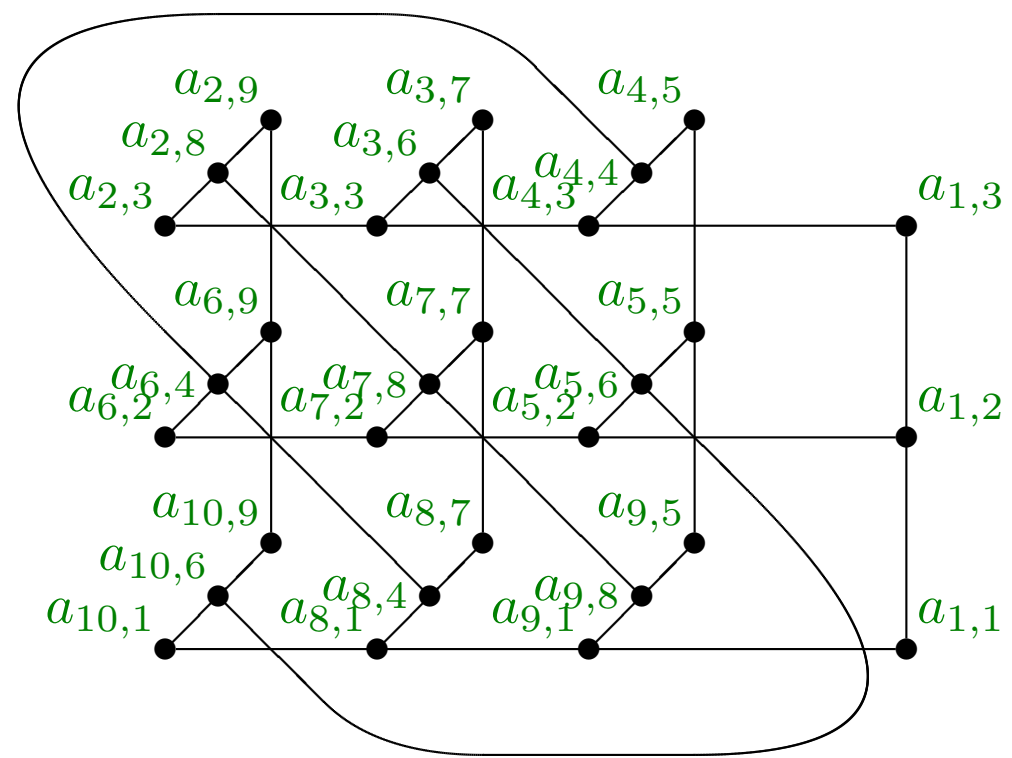




Smaller example with empty state space [Greechie]:



Even smaller example with empty state space [R. Mayet]:



This is the smallest example with empty state space obtained by this technique and it is not unique [MN 08]; it has 19 blocks.

OMLs with  $\leq 5$  blocks admit states [Riečanová 07].

## Bell inequalities

$$s(a) + s(b) - s(a \wedge b) \leq 1$$

$$0 \geq s(a \wedge b) + s(b \wedge c) + s(c \wedge d) - s(a \wedge d) - s(b) - s(c)$$

$$s(a) + s(b) + s(c) - s(a \wedge b) - s(a \wedge c) - s(b \wedge c) \leq 1$$

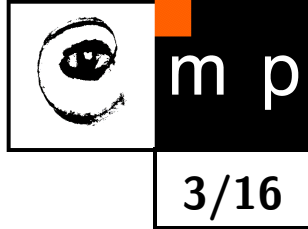
$$s(a \wedge b) + s(b \wedge c) + s(c \wedge d) - s(a \wedge d) - s(b) - s(c) \geq -1$$

The first is equivalent to the **valuation property**:

$$s(a \wedge b) + s(a \vee b) = s(a) + s(b)$$

If the OML is not a Boolean algebra and admits a rich set of states, all Bell inequalities are violated.

# Crucial example of a quantum structure: Hilbert lattice



$H$  ... a separable Hilbert space (real or complex)

$L(H)$  ... the set of all closed subspaces of  $H$  (equivalently, all projectors of  $H$ )

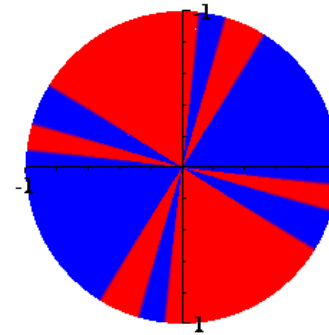
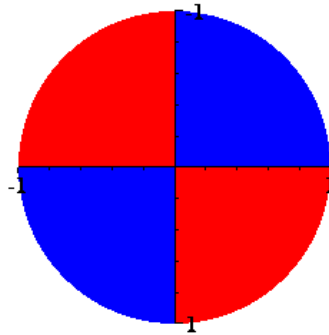
$$\begin{aligned} \mathbf{0} &= \{0\}, \\ \mathbf{1} &= H, \\ A \leq B &\iff A \subseteq B, \\ A \wedge B &= A \cap B, \\ A' &= \{x \in H \mid \forall y \in A : y \perp x\}, \\ A \vee B &= \text{Lin}(A \cup B), \end{aligned}$$

where  $\text{Lin}$  denotes the closed linear hull



The only restriction of states for  $\dim P = 1$ :  $s(P') = 1 - s(P)$

Many two-valued states = colourings of non-zero vectors by two colors (blue, red) such that each orthogonal basis contains exactly one red vector



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1. For  $q \in H$ ,  $\|q\| = 1$ , define a **vector state**

$$s_q(\text{Lin}(\{y_1, \dots, y_n\})) = \sum_{i=1}^n (q \cdot y_i)^2 = \sum_{i=1}^n \cos^2 \angle(q, y_i)$$

for any orthonormal basis  $(y_1, \dots, y_n)$  of  $H$

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**Corollary:**  $s_q(q) = 1$

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**Corollary:**  $s_q(q) = 1$

2. **Mixture** of vector states

$$s(P) = \sum_i c_i s_{q_i}(P), \text{ where } c_i > 0, \sum_i c_i = 1 .$$

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3. What else?

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2. **Mixture** of vector states  $s(P) = \sum_i c_i s_{q_i}(P)$ , where  $c_i > 0$ ,  $\sum_i c_i = 1$ .

3. What else?

**Nothing!**

**Gleason's Theorem** [Gleason 57]: For  $\dim H \geq 3$ , all states are mixtures of vector states.

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**Crucial case:**  $H = \mathbb{R}^3$  (simplified proof by [Cooke, Keane, Moran 85]).

**Corollary 1:** The restriction of a state to 1D subspaces is continuous (proved by [von Neumann 1932] even for  $\mathbb{R}^2$ , error found by [Hermann 1935],

**Corollary 2:** A finitely-valued state is constant on 1D subspaces, i.e., it is a dimension function.

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**Crucial corollary:** There is no two-valued state (=hidden variable) (answer to a question by [Einstein, Podolsky, Rosen 35], simplified proofs by [Bell 64, Bell 66], [Kochen, Specker 67], ... , [Peres 95]).

1	2
3	4
5	6
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17	18
19	20
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25	26
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29	30



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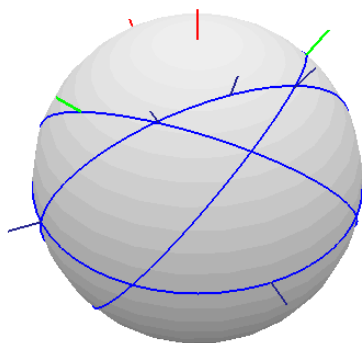
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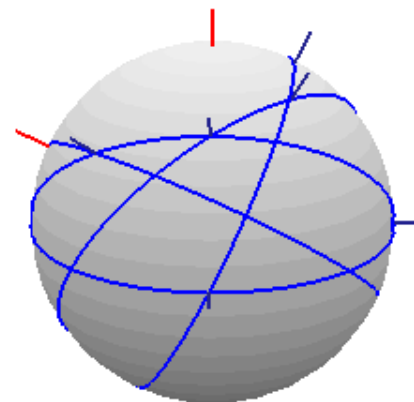
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# Constructions proving Geometrical Lemmas

BGL



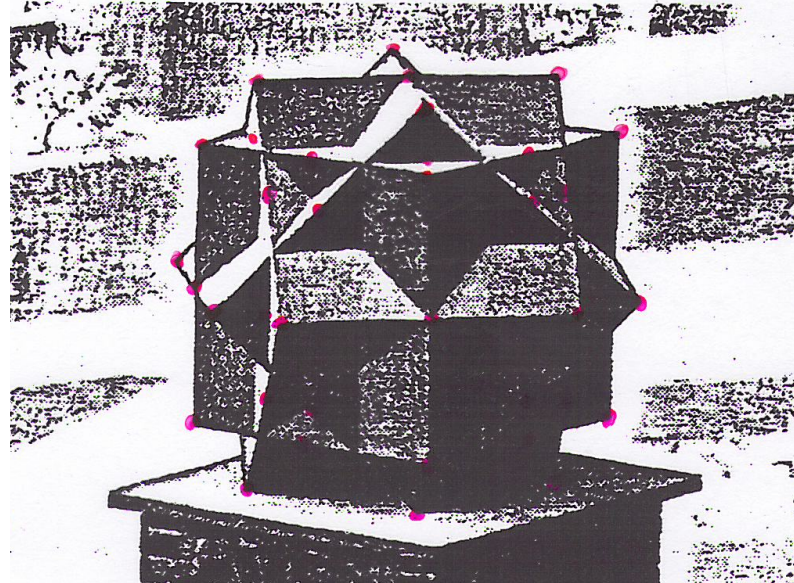
MGL





It is possible to find a finite set of vectors whose orthogonality relations exclude the possibility of a two-valued state.

The smallest example known uses 31 vectors, the following uses 33 vectors:



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**Theorem:** [Cabello] There is no two-valued state on  $\mathcal{L}(\mathbb{R}^4)$ .

Take 36 vectors in  $\mathbb{R}^4$  ( $\bar{1}$  denotes  $-1$ ):

1000	1000	0100	1111	1111	111 $\bar{1}$	11 $\bar{1}\bar{1}$	111 $\bar{1}$	11 $\bar{1}1$
0100	0010	0010	11 $\bar{1}\bar{1}$	1 $\bar{1}1\bar{1}$	11 $\bar{1}1$	1 $\bar{1}1\bar{1}$	1 $\bar{1}11$	1 $\bar{1}11$
0011	0101	1001	1 $\bar{1}00$	10 $\bar{1}0$	1 $\bar{1}00$	1001	10 $\bar{1}0$	100 $\bar{1}$
001 $\bar{1}$	010 $\bar{1}$	100 $\bar{1}$	001 $\bar{1}$	010 $\bar{1}$	0011	0110	0101	0110

Each of the 9 column represents an orthogonal basis of  $\mathbb{R}^4$  and each vector **occurs twice**. The number of vectors of unit state in this table must be both even and odd (9)—a contradiction.

1	2
3	4
5	6
7	8
9	10
11	12
13	14
15	16
17	18
19	20
21	22
23	24
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27	28
29	30



## Corollaries for **group-valued** states

**Theorem:** There is no nontrivial  $\mathbb{Z}_2$ -valued state,  $s$ , on  $\mathcal{L}(R^4)$  which satisfies  $s(\mathbf{1}) = 1$ .

**Theorem:** If  $n \geq 4$ , then there is no nontrivial  $\mathbb{Z}_2$ -valued state on  $\mathcal{L}(R^n)$ .

For  $n \geq 5$  it follows from the above construction.

The refinement for  $n = 4$  is due to [Harding, Jager, Smith]

**Open problem:** Are there nontrivial  $\mathbb{Z}_2$ -valued states on  $\mathcal{L}(R^3)$ ?

## References

- [Bell 64] Bell, J.S.: On the Einstein–Podolsky–Rosen paradox. *Physics* **1** (1964), 195–200.
- [Bell 66] Bell, J.S.: On the problem of hidden variables in quantum theory. *Rev. Mod. Phys.* **38** (1966), 447–452.
- [Birkhoff, vonNeumann 36] Birkhoff, G., von Neumann, J.: The logic of quantum mechanics. *Ann. Math.* **37** (1936), 823–843.
- [Cooke, Keane, Moran 85] Cooke, R., Keane, M., Moran, W.: An elementary proof of Gleason’s theorem. *Math. Proc. Cambridge Philos. Soc.* **98** (1985), 117–128.
- [Dvurečenskij 93] Dvurečenskij, A.: *Gleason’s Theorem and Its Applications*. Kluwer, Dordrecht/Boston/London & Ister Sci., Bratislava, 1993.
- [Einstein, Podolsky, Rosen 35] Einstein, A., Podolsky, B., Rosen, N.: Can quantum-mechanical description of physical reality be considered complete? *Phys. Rev.* **47** (1935), 777–780.

## References

- [Gleason 57] Gleason, A.M.: Measures on the closed subspaces of a Hilbert space. *J. Math. Mech.* **6** (1957), 885–893.
- [Greechie 71] Greechie, R.J.: Orthomodular lattices admitting no states. *J. Comb. Theory* **10** (1971), 119–132.
- [Gudder 66] Gudder, S.P.: Uniqueness and existence properties of bounded observables. *Pacific J. Math.* **19** (1966), 81–93.
- [Harding, Jager, Smith] Harding, J., Jager, K., Smith, D.: Group-valued measures on the lattice of closed subspaces of a Hilbert space. *Internat. J. Theoret. Phys.* **44**, no. 5 (2005), 539–548.
- [MN 95] Navara, M.: Uniqueness of bounded observables. *Ann. Inst. H. Poincaré — Theor. Physics* **63** (1995), no. 2, 155–176.
- [Kalmbach 83] Kalmbach, G.: *Orthomodular Lattices*. Academic Press, London, 1983.
- [MN 92] Navara, M.: Descriptions of state spaces of orthomodular lattices. *Math. Bohem.* **117** (1992), 305–313.

[MN 00] Navara, M.: State spaces of orthomodular structures. *Rend. Istit. Mat. Trieste* **31** (2000), Suppl. 1, 143–201.

[MN 04] Navara, M.: Piron's and Bell's geometrical lemmas. *Internat. J. Theoret. Phys.* **43** (2004), No. 7, 1587–1594.

[MN 08] Navara, M.: Small quantum structures with small state spaces. *Internat. J. Theoret. Phys.* **47** (2008), No. 1, 36–43.

[Pták, MN 04] Navara, M., Pták, P.: For  $n \geq 5$  there is no nontrivial  $Z_2$ -measure on  $L(R^n)$ . *Internat. J. Theoret. Phys.* **43** (2004), No. 7, 1595–1598.

[Pták, Pulmannová 91] Pták, P., Pulmannová, S.: *Orthomodular Structures as Quantum Logics*. Kluwer, Dordrecht/Boston/London, 1991.

[Riečanová 07] Riečanová, Z.: The existence of states on every Archimedean atomic lattice effect algebra with at most five blocks. Preprint, 2007.

[Shultz 74] Shultz, F.W.: A characterization of state spaces of orthomodular lattices. *J. Comb. Theory A* **17** (1974), 317–328.

# Details of pictures

